#### **Chiara Saffirio**





Geneva, YRS 2021

Boltzmann (1872) and Maxwell (1867) attempt at a realistic description of rarefied gases





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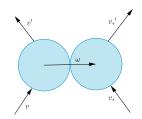
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$$Q(f,f) = \int_{\mathbb{R}^3} dv_* \int_{S^2} d\omega B(\omega, v - v_*)$$

$$\times \{f(t, x, v_*')f(t, x, v_*') - f(t, x, v_*)f(t, x, v_*)\}$$



### **Conservation laws and H-Theorem**

Mass, Momentum, Energy

$$\iint_{\mathbb{R}^3 \times \mathbb{R}^3} \varphi(v) f(t, x, v) dx dv = \iint_{\mathbb{R}^3 \times \mathbb{R}^3} \varphi(v) f_0(x, v) dx dv$$

f solution to the Boltzmann eq. with initial datum  $f_0$  and  $\varphi(v) = 1$ ,  $v_i$ ,  $v^2$ .

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### Theorem (H Theorem, Boltzmann '72)

If f(t) is a regular enough solution to the Boltzmann equation, then

$$H(t) \leq H(0)$$

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#### Global well-posedness?

On the one hand:

$$(\partial_t + v \cdot \nabla_x) f(t, x, v)$$

$$= \int_{\mathbb{R}^3} dv_* \int_{S^2} d\omega B(\omega, v - v_*) \{ f(t, x, v_*') f(t, x, v') - f(t, x, v_*) f(t, x, v) \}$$

looks like

$$\partial_t f(t) \sim f(t)^2 \implies \text{only local in time!}$$

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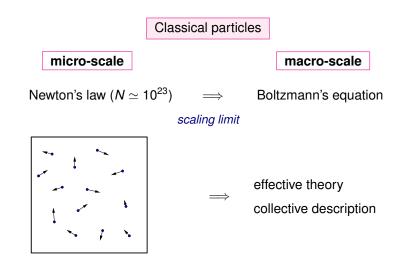
#### Global well-posedness?

On the other hand:

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there might be cancellations!



#### Newton: time reversible dynamics

$$\begin{cases} \frac{d}{dt}x_i(t) = v_i(t), \\ \frac{d}{dt}v_i(t) = 0, \\ i = 1, \dots, N \end{cases}$$
 + boundary conditions

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#### Liouville equation:

$$\partial_t f_N + \sum_{i=1}^N v_i \cdot \nabla_{x_i} f_N = 0 + \text{b.c.}$$

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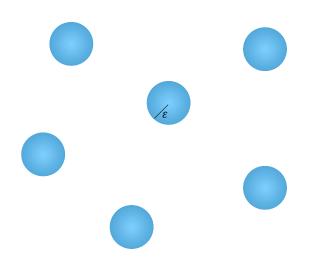
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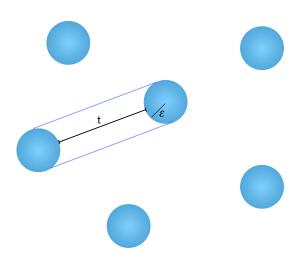
### *j*-particle marginal:

$$f_N^{(j)}(t, x_1, v_1, \dots, x_j, v_j) = \int f_N(t, x_1, v_1, \dots, x_N, v_N) dx_{j+1} dv_{j+1} \dots dx_N dv_N$$

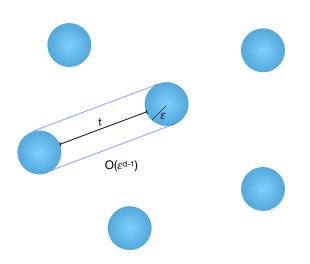
The Boltzmann-Grad limit: N particles of radius  $\varepsilon$ ,  $N \to \infty$  and  $\varepsilon \to 0$ 



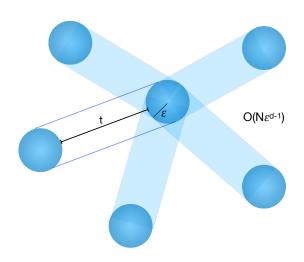
The Boltzmann-Grad limit: N particles of radius  $\varepsilon$ ,  $N \to \infty$  and  $\varepsilon \to 0$ 



The Boltzmann-Grad limit: N particles of radius  $\varepsilon$ , low density regime



The Boltzmann-Grad limit:  $N \to \infty$  with the constraint  $N\varepsilon^{d-1} = O(1)$ 



$$\partial_t f_N + \sum_{i=1}^N v_i \cdot \nabla_{x_i} f_N = 0 + \text{b.c.}$$

and consider the first marginal  $f_N^{(1)}$ .

$$\begin{split} (\partial_t + v \cdot \nabla_x) f_N^{(1)}(t,x,v) &= (N-1)\varepsilon^2 \int_{\mathbb{R}^3} \int_{S^2} B(\omega,v-v_*) \\ &\times \{ f_N^{(2)}(t,x-\varepsilon\omega,v_*',x,v') - f_N^{(2)}(t,x+\varepsilon\omega,v_*,x,v) \} \, d\omega \, dv_* \end{split}$$

to be compared with

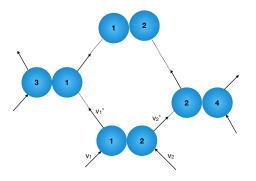
$$(\partial_t + v \cdot \nabla_x) f(t, x, v) = \int_{\mathbb{R}^3} \int_{S^2} B(\omega, v - v_*) \times \{ f(t, x, v_*') f(t, x, v') - f(t, x, v_*) f(t, x, v) \} d\omega dv_*$$

Propagation of chaos

$$f_N^{(2)}(0) \sim f_0^{\otimes 2} \quad \Longrightarrow \quad f_N^{(2)}(t) \sim f(t)^{\otimes 2}$$

where f is a solution of the Boltzmann equation with initial datum  $f_0$ .

Accurate study of pathological configurations



- Lanford (1975): hard spheres, short times
- Gallagher, Saint-Raymond, Texier (2013): quantitative analysis

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#### Class of interactions:

- Short range potentials:
   Gallagher, Saint-Raymond, Texier (2013), Pulvirenti, C.S., Simonella (2014)
- \* Triple interactions: Ampatzoglou, Pavlovic (2019, 2020)
  - Long range potentials: (?

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- \* Long range potentials: ?

#### Time of validity:

- \* Near the vacuum: Illner, Pulvirenti (1986)
- Linear and linearized setting:
   Bodineau, Gallagher, Saint-Raymond (2016, 2017), + Simonella (2020)
- \* Nonlinear setting: (?) (related to the global existence for the PDE)

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...and many other open problems (boundaries, boundary layers, molecular interactions, ...)
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# Plenty of work to be done!